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The first observation in an organic medium and DFT calculation of the TMM radical anion generated via a single electron reduction of a methylenecyclopropane

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Abstract— γ -Irradiation of 2,2-diphenyl-1-methylenecyclopropane (3) in a degassed 2-methyltetrahydrofuran glassy matrix at 77 K gave an intense UV/vis absorption band with λ_{ab} at 496 nm. This result and calculations based on density functional theory for its radical anion 3⁻ and the corresponding trimethylenemethane radical anion (2⁻) strongly suggest that single electron reduction of 3 followed by ready ring opening affords 2⁻, whose molecular geometry is largely twisted ($\theta = 45.5^{\circ}$), and the negative charge and spin are localized mainly in the diphenyl methyl and allyl moieties, respectively. © 2006 Elsevier Ltd. All rights reserved.

Since Dowd first generated the parent trimethylenemethane (TMM, 1" in Chart 1),¹ various neutral TMM derivatives have been generated and studied from a variety of different perspectives.² Miyashi and Roth³ were the first to generate a TMM radical cation derivative, $(2^{\cdot+})$, via a photoinduced electron-transfer reaction using 2,2-diphenyl-1-methylenecyclopropane (3), and Ikeda and Miyashi⁴ confirmed its reaction mechanism, mole-



Chart 1. A list of TMM derivatives and structurally related compounds.

cular geometry (largely twisted, $\theta \approx 44^{\circ}$, vide infra), and electronic structure (considerably divided) in spectroscopic, calorimetric, and semi-empirical calculation studies using its bis(4-methoxyphenyl) derivative.⁵

The generation of $2^{\cdot+}$ from 3 is responsible for the C-2– C-3 bond (Chart 1) cleavage of $3^{\cdot+}$ induced by ejection of an electron from the highest occupied molecular orbital (HOMO) of 3 with bonding character. In other words, lowering the bond order of $3^{\cdot+}$ holds the key to produce $2^{\cdot+}$. An alternative method of lowering the bond order is to inject an electron into the lowest unoccupied molecular orbital (LUMO) with antibonding character. Namely, 3 should have the potential to generate the corresponding TMM radical anion ($2^{\cdot-}$) via a single electron reduction.

This concept is of interest because there are few studies of the parent TMM radical anion (1⁻⁻) or its derivatives. As far as we know, the sole example is the nonreductive generation of 1⁻⁻ using the reaction of 3-(trimethylsilyl)-2-(trimethylsilylmethyl)propene with fluoride ion and molecular fluorine in the gas phase by Squires, Lineberger, and co-workers,⁶ and there is no example of their generation or observation in an organic medium. Therefore, to test this concept, we conducted γ -irradiation of **3** in a 2-methyltetrahydrofuran (MTHF) glassy matrix at 77 K, which is a general method for generating radical anions. Here, we report the first isolation of TMM radical anion derivatives, that is **2**⁻⁻, in an MTHF matrix,

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together with its molecular geometry and electronic structure evaluated using density functional theory (DFT) calculations at the UB3LYP/cc-pVDZ level of theory.

A degassed MTHF glassy matrix of 3 (5 mM) was γ irradiated with a ⁶⁰Co source at 77 K.⁷ As shown in Figure 1, an intense absorption band appeared with λ_{ab} at 496 nm. Two assignments for the 496-nm band are possible: the precursor radical anion 3^{-} and the TMM radical anion 2^{-} generated via 3^{-} (Scheme 1). For comparison, 1,1-diphenylcyclopropane (4, 5 mM) and 2,2-diphenylpropane (5, 5 mM) were similarly γ -irradiated in degassed MTHF glassy matrices at 77 K. As a result, neither 4^{•-} nor 5^{•-} had any significant absorption bands in the 300-600 nm region.^{7,8} Conversely, the 1,1diphenylethyl anion (6^{-}) has an absorption band with $\lambda_{ab} = 477 \text{ nm}$ in cyclohexylamine,⁹ which is the sole study on the absorption band of 6^- . These findings strongly suggest that the observed absorption band $(\lambda_{ab} = 496 \text{ nm}, \text{ Fig. 1})$ is assigned to the TMM radical anion 2^{-} rather than 3^{-} and that its negative charge and spin are considerably localized in the diphenylmethyl (subunit I) and allyl (subunit II) moieties (see 2⁻⁻ in Chart 1), respectively (vide infra). However, some 3^{-} may still remain in the matrix though there is no spectroscopic evidence for 3^{.-}, unfortunately. Remember that 3^{-} is not expected to have an absorption band in the 300–600 nm region, similar to 4^{-} and 5^{-} .

To gain insight into the ring opening process from 3^{-1} to 2^{-1} , the molecular geometry and electronic structure, we performed DFT calculations at the UB3LYP/cc-pVDZ level.^{10,13} Figure 2 shows the C-2–C-3 bond length ($d_{C-2-C-3}$, see Chart 1)-dependent change in the relative potential energy (ΔE) for ring opening. In Figure 2, the leftmost and rightmost points with $d_{C-2-C-3} = 1.58$ and 2.52 Å are for the optimized 3^{-1} and 2^{-1} ,



Figure 1. UV/vis absorption spectrum observed after γ -irradiation of **3** (5 mM) in an MTHF glassy matrix at 77 K.



Scheme 1. Possible generation of the TMM radical anion 2^{-} by γ -irradiation of **3** in an MTHF glassy matrix at 77 K.



Figure 2. Bond length ($d_{C-2-C-3}$)-dependent changes in the relative potential energy (ΔE) in the course of ring opening from 3⁻⁻ to 2⁻⁻ calculated using UB3LYP/cc-pVDZ.

respectively. The ΔE value increased to reach a maximum at $d_{C^{-2}-C^{-3}} \approx 1.8$ Å, and then decreased monotonically, as $d_{C^{-2}-C^{-3}}$ increased from 1.58 to 2.52 Å. Estimating from ΔE , the ring opening of **3**⁻ to give **2**⁻ proceeds readily with a relatively small elongation barrier (ca. 2 kcal mol⁻¹) and the release of a large amount of energy (ca. 28 kcal mol⁻¹). This finding rationalizes our assignment of the spectrum observed in the γ -irradiated MTHF glassy matrix of **3** (Fig. 1, vide supra). Again, note that some **3**⁻ may still survive the calculated low elongation barrier at 77 K.

Figure 3 shows the optimized molecular geometry of 2⁻⁻. The dihedral angles, θ and θ' , of C-3–C-1–C-2–C-5 and C-3–C-1–C-2–C-7 (Fig. 3d) were optimized to be +45.5° (Fig. 4) and –134.5°, respectively, while similar angles, ω and ω' , of C-1–C-2–C-5–C-6 and C-1–C-2–C-7–C-8 were both calculated to be +25.3°. The



Figure 3. Molecular geometry of 2^{-} optimized using UB3LYP/ccpVDZ: (a) top, (b) front, and (c) side views. (d) Atom notation and definition of the dihedral angles $(\theta, \theta', \omega, \text{ and } \omega')$ of 2^{-} .



Figure 4. Potential energy (ΔE) curve of **2**⁻⁻ along the dihedral angle (θ) calculated using UB3LYP/cc-pVDZ.

sum of calculated internal angles $\angle C-1-C-2-C-5 =$ 118.6°, \angle C-5–C-2–C-7 = 122.8°, and \angle C-7–C-2–C-1 = 118.6°, was 360°, suggesting the sp²-like hybridization of C-2 in 2⁻⁻. In fact, it was suggested that the two σ bonds of the C-2 carbon of 2^{-} are hybridized with (s,p) = (34.20%, 65.74%) in the directions of C-5 and C-7, whereas in the direction of C-1, the residual σ bond of C-2 carbon takes an (s,p) = (31.52%, and 68.42%)hybridization.¹⁴ Note that C-2 of 2^{-1} is sp²-like in hybridization while the methyl anion, CH_3^- , has an sp³ hybridized carbon. This phenomenon is due to both electronic and geometric factors. The former is a conjugation of the negative charge to the two phenyl groups of 2^{-} , while the latter is the easing of potential steric hindrance between the two phenyl groups on C-2 with a hypothetical sp³ hybridization.

These findings, especially θ and θ' , suggest that the ground state of 2^{-1} is largely twisted. The optimized C-1–C-2 bond length ($d_{C-1-C-2}$), 1.49 Å, is slightly longer than the C–C bond length ($d_{C-C} = 1.47$ Å) of the orthogonal ethylene.¹⁵ Therefore, the C-1–C-2 bond of 2^{-1} possesses a normal C–C single bond character. Furthermore, judging from the ΔE of 2^{-1} along the θ coordinate (Fig. 4), the rotation barrier around the C-1–C-2 bond is lesser than 4 kcal mol⁻¹ at $\theta \approx 0^{\circ}$, which corresponds to that of the normal C–C single bond (3–6 kcal mol⁻¹), indicating that the C-1–C-2 bond of 2^{-1} does not have the typical nature of a double bond.

Finally, we examined the electronic structure using the sum of the partial spin (ρ) and charge (q) density, $\sum \rho$ and $\sum q$, respectively, computed with DFT calculations, as shown in Figure 5.¹⁶ The values, $\sum \rho = -0.025$ and $\sum q = -0.782$ for subunit I, and $\sum \rho = +1.025$ and $\sum q = -0.218$ for subunit II suggest that the negative



Figure 5. The sum of the partial spin (ρ) and charge (q) density of subunits I and II in 2^{-} , respectively.

charge and spin are mainly localized in subunits I and II, respectively. These calculation results are in accord with the conclusion based on absorption spectroscopy.

In conclusion, we first observed the TMM radical anion derivative, that is 2^{-} , generated in organic medium using γ -irradiation of 3 in an MTHF glassy matrix at 77 K. The C-2–C-3 bond cleavage of 3^{-1} results from lowering the bond order triggered by an electron injection to the LUMO with antibonding character. The DFT calculation at the UB3LYP/cc-pVDZ level suggests that 2^{-} has a largely twisted molecular geometry and a considerably localized electronic structure. Subunits I and II are largely twisted with each other around the C-1-C-2 bond. The negative charge and spin are mainly distributed to subunits I and II, respectively. These phenomena are also similar to the case of 2^{+} : the corresponding subunits I and II are largely twisted with each other around the C-1-C-2 bond, and the positive charge and spin are mainly distributed to the corresponding subunits I and II, respectively.¹⁷ This work is significant from the perspective of providing a new system of radical anion rearrangement, because such reaction systems are not documented very often,18 as compared to those of radical cation rearrangements. Further studies (e.g., comparison with the parent 1^{-}) are now in progress, and will be published elsewhere.

Acknowledgments

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Supplementary data

The DFT calculation result for 2^{-} and the absorption spectra of the MTHF matrices of 4^{-} and 5^{-} obtained after γ -irradiation at 77 K. Supplementary data associated with this article can be found, in the online version, at doi:10.1016/j.tetlet.2006.01.024.

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